**Double-end pumped efficient MHz self-mode-locked Pr:YLF green and red lasers**

**Saiyu Luo, Bin Xu\*, Huiying Xu, Zhiping Cai1,\***

*Department of Electronic Engineering, Xiamen University, Xiamen 361005, China*

*\*E-mail:* [xubin@xmu.edu.cn](mailto:xubin@xmu.edu.cn), zpcai@xmu.edu.cn

**Abstract:** We demonstrate an efficient tens and hundreds of MHz self-mode-locked green and red laser in a Pr3+:YLF4 crystal with V-type and Z-type cavity using a double-end-diode-pumped structure without the need of any additional components. Experimental results reveal that the laser system can be characterized in stable mode-locked operations. With an absorbed pump power of ~2.8 W, average output powers of more than 0.68 W for 522 nm green laser or more than 1.44 W for 639 nm red laser were obtained.

**Keywords:** Solid state laser, Blue diode-double-end pump, Pr:YLF crystal,visible laser source.

**1. Introduction**

The third-order nonlinear optical responses are closely related to the stimulated Raman

scattering (SRS) process and the Kerr-lensing effect [].

Gaponenko et al. obtained a pulse width of 18 ps (FWHM) with SESAM mode-locking [SESAM mode locked red praseodymium laser].

**2. Experimental Set-Up**

Figure 1 depicts the experimental setup for the self-mode-locked Pr:YLF red laser with a repetition rate of ~85 MHz using a double-end pumping scheme. The cavity configuration is a Z-type resonator with a total cavity length of ~1.8 m. The active medium is an a-cut 0.2 at. % Pr3+:YLF4 crystal with a length of ~8 mm. Both end surfaces of the Pr3+:YLF4 crystal were plano and uncoated. The laser crystal was wrapped with indium foil and mounted in a water-cooled copper holder. The water temperature was maintained around 18°C to eliminate thermal effect. The laser crystal was placed close to IM1, which is the input mirror for the pump on the left side. IM2 is the input mirror for the pump on the right side and was tilted to make the cavity longer. Both IM1 and IM2 had high transmission (>95%) at ~444 nm and high reflection (>99.8%) at 639 nm.

Each of the two pump sources was a ~1.8 W, ~444 nm laser diode with achromatic and collimation system, which is commercially available. Focusing lenses with 50 mm focal length was used to inject the pump beam into the laser crystal with an average pump size of approximately 60 μm, which is measured by using Spiricon M2-200.

HR1~HR3 were concave mirrors with radius of curvature of 300 nm and had high reflection (>99.7%) at 639 nm, which could provide positive feedback for 639 nm lasing. The output coupler (OC) was a coated plane mirror with a transmission of 3.5% at 639 nm. The distance between IM1 and IM2 plus the distance between IM2 and HR1 equals ~300 mm. The distance between HR1 and HR2 and the distance between HR2 and HR3 were both around 600 mm. The distance between HR3 and OC was ~300 mm. So the total cavity length was ~1.8 m.



Fig. 1. Schematic of the Z-type self-mode-locked Pr3+:YLF4 red laser.

Moreover, for 639 nm self-mode-locked lasing operation, experiments with a V-type cavity and concave OC with a radius of curvature of 300 mm was conducted, making the cavity length ~0.3 m. The same V-type schematic using different IMs and OC for 522 nm lasing was also used to generate green self-mode-locked laser. For the 522 nm self-mode-locked laser, the two input mirrors IM1 and IM2 had a high transmission (> 95%) for the pump (~444 nm), high reflection (> 99.9%) at 523 nm to support lasing and high transmission (> 60%) around 607 and 639 nm to suppress the high gain emissions at those wavelengths. The concave output coupler had a radius of curvature of 300 mm and transmission of 1.9% at 522 nm.

**3. Experimental results and discussions**

For the Z-type cavity with a total cavity length of ~1.8 m shown in Figure 1 operating at 639 nm, the optical cavity length was firstly set to be approximately 1.8 m, corresponding to a FSR of 83.3 MHz. Than the cavity alignment was optimized for generating the maximum average output power. The time trace of the output radiation revealed the laser to be operating in the self-mode-locked state. Thanks to the double-end pumping scheme, more than 1.44 W output power was extracted with a total laser slop efficiency of 57.3% with respect to the absorbed pump power. In the experiment, the left and right pump LDs were turned on in turn. The lasing performance for ~1.8 m Z-type cavity operated in self-mode-locked state at the wavelength of 639 nm was shown by the red dots and line depicted in Figure 2.

For the green self-mode-locked Pr3+:YLF4 laser operating at the wavelength of 522 nm, which adopted a ~0.3 m V-type cavity using OC with a transmission of 1.9%, owing to the double-end pumping scheme, up to ~0.68 W average output power at mode-locked state was obtained, with a total laser slop efficiency of 29.9% with respect to the absorbed pump power as shown by the green dots and line depicted in Figure 2.

The reason why we did not manage to extend the cavity length of the green self-mode-locked laser might result from the fact that the green laser is more susceptible to the air conditions, which might limit the transmitting distance in the atmosphere, especially when considering intra-cavity transmission.

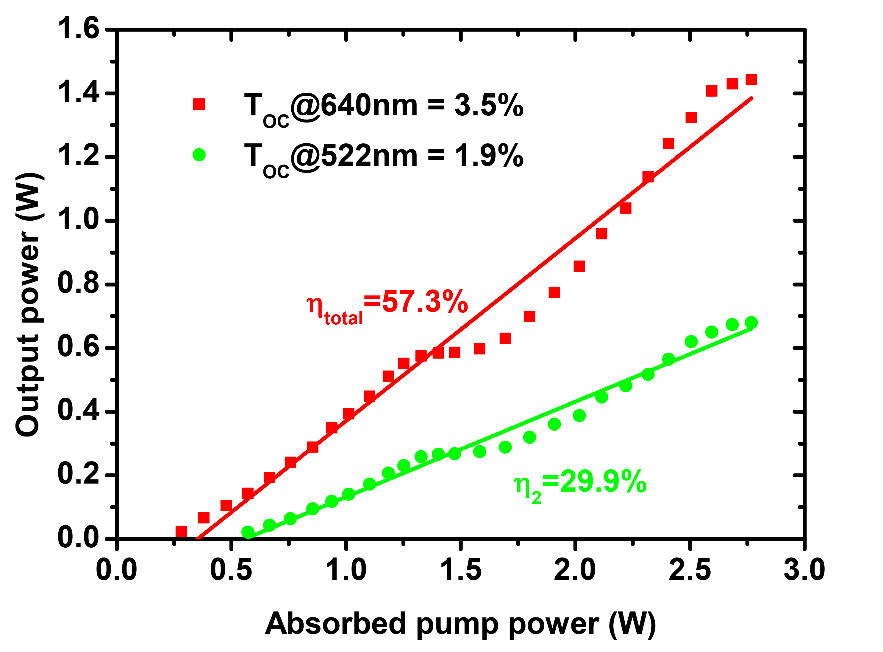


Fig. 2. Self-mode-locked laser performance of the ~1.8 m cavity 639 nm (red) and ~0.3 m 522 nm (green) laser operations.

It is worth noting that the flat points at the absorbed pump power of 1.4~1.6 W shown in Figure 2 were caused by the introduction of double-end pumping scheme and the way we record the input-output relation, which is by first turning one pump to maximum and then another. Another notable feature for the laser performance is that the double-end pumping scheme could lead to a larger slop efficiency compared with individual pumping, which might belong to the increased overlapping efficiency that the pump interacts with laser cavity modes when two pump beams were injected from both ends.

The optical spectrum of the ~1.8 m cavity 639 nm laser operated in self-mode-locked state was shown in Figure 4. The spectrum was registered by Hewlett Packard 8560E Series Optical Spectrum Analyzer with a resolution of 0.08 nm under maximum pump power and meanwhile, a photo detector and digital oscilloscope was used to monitor the laser output to ensure that the laser was operated in mode-locked state. As shown in Figure 4, the center wavelength was 639.4 nm, and the spectral width (FWHM) was 0.17 nm, which corresponded to a transform-limited pulse width of 2.6 ps for Gaussian-shaped pulses. In fact, the measured spectrum shows some internal structures.

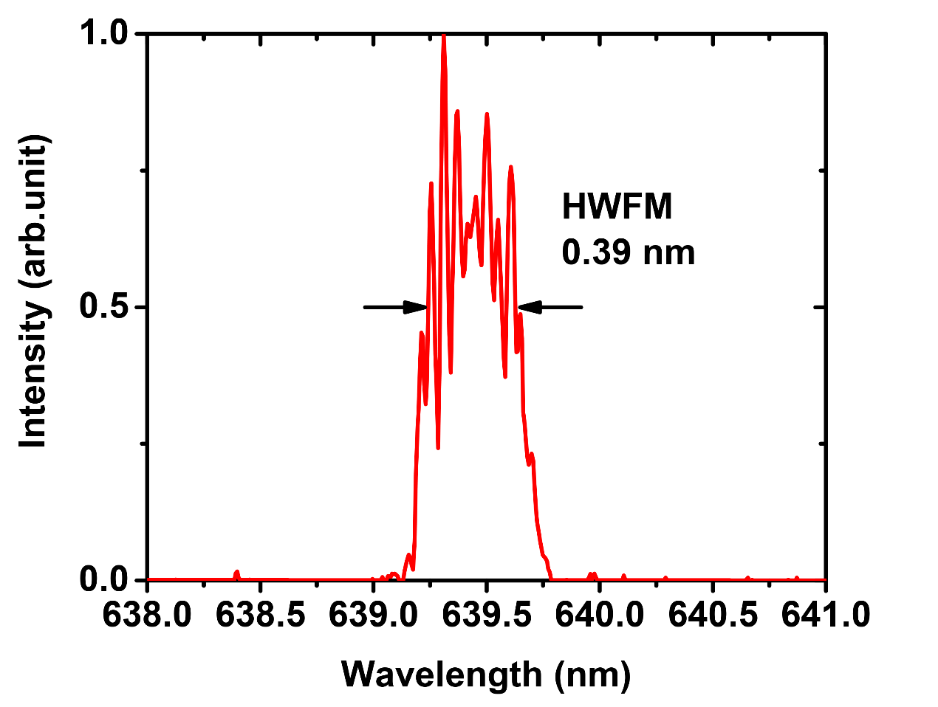


Fig. 3. Optical spectrum of the self-mode-locked Pr3+:YLF4 laser for 639 nm at cavity length of ~1.8 m measured at the maximum pump power.

The mode-locked pulses trains were detected by a free space high sensitivity PIN photo detector unit (Menlo Systems, Inc. FPD310-FV with rise time 0.7 ns), whose output signal was connected to a digital mixed signal oscilloscope (Tektronix MSO 3054) with 500 MHz electrical bandwidth and a sampling interval of 0.4 ns.

Figure 4 shows the pulse train results measured from the output of one of the high reflection mirrors (HR1~HR3) or IM2 at the same time the maximum emitting laser powers from output coupler (OC) were measured. This has the advantage of both ensuring that the lasers were in mode-lock state while optimizing output power, and protecting the photo detector by letting few energies hit on it.

Figure 4 shows the pulse trains for the self-mode-locking operated at the wavelength of 522 nm (Figure 4(a), Figure 4(d)) and 639 nm (Figure 4(b), Figure 4(e)) with a cavity length of ~0.3 m, and the self-mode-locking pulse trains at the wavelength of 639 nm with a cavity length of ~1.8 m (Figure 4(c), Figure4(f)), respectively, on two different time scales.

The figures on the left side of Figure 4 are the ones with time span of 5 ns, demonstrating mode-locked pulses, and the figures on the right side of Figure 4 are the ones with time span of 5 μs, demonstrating the amplitude stability. It can be seen that the pulse trains display full modulation, and the complete mode locking is achieved. In fact, the stable mode-lock oscillation could be maintained for few hours at the maximum output power.

Unfortunately, the sampling rates of our photo detector and digital oscilloscope were not sufficient enough to detect the details of the pulse trains and the single pulse traces could neither be depicted, because the auto-correlation function analyzer in our lab was not available in the visible spectral range.



Fig. 4. Output pulse trains of the self-mode-locked lasers operating at the wavelength of 522 nm and 639 nm with a cavity length of ~0.3 m ((a), (b), (d), (e)) and ~1.8 m ((c), (f)) in time span of 5 ns ((a), (b), (c)) and 5 μs ((d), (e), (f)).

Note that once the pump power reaches the lasing threshold, the laser system instantaneously steps into a stable mode-locked operation without any mechanical perturbation. The locking mechanism is presumed to be the Kerr effect. However, the laser system has high stability over day-long operation and is insensitive to mechanical vibrations and air current. As a result, some auxiliary mechanism seems to exist in the locking process. Bai et al. [Novel self-mode-locking mechanism in narrow-band lasers] proposed a novel self-mode-locking mechanism in narrowband lasers based on the analysis of the gain-line splitting induced by an intra-cavity laser field. Although the present experimental results are fairly consistent with this mechanism, further identification is still needed.

remains an open question

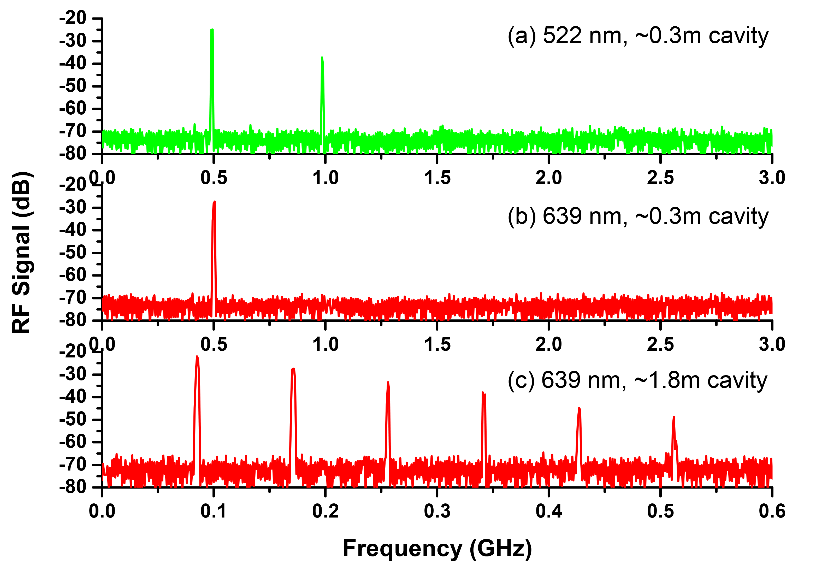


Fig. 5. Power spectrum of the self-mode-locked lasers operating at the wavelength of 522 nm and 639 nm with a cavity length of ~0.3 m ((a), (b)) and ~1.8 m ((c))

The power spectrum of the mode-locking outputs were detected by the same free space high sensitivity PIN photo detector unit as the one used for mode-locked pulse train measurement, whose output signal was transferred to an rf spectrum analyzer (GωINSTEK, GSP-930) with a bandwidth of 3.0 GHz. The results shown in Figure 5 were also recorded through one of the three high reflection mirrors (HR1~HR3) or IM2 when maximum output powers were achieved.

As is depicted in Figure 5(a), repetition rate of 520 MHz was registered for self-mode-locked laser operating at the wavelength of 522 nm using V-type cavity and OC with radius of curvature of 300 mm, indicating an effective cavity length of 288 mm when the cavity was optimized for the maximum output power.

On the other hand, by introducing the same V-type cavity and the same radius of curvature of 300 mm for the OC, but switching cavity mirrors for the 639 nm laser emitting, repetition rate of 535 MHz was recorded for the self-mode-locked laser operating at the wavelength of 639 nm as shown in Figure 5(b), deducing an effective cavity length of 280 mm when the maximum output power was achieved.

The reason why the optimized cavity length for the 522 nm laser was longer than the 639 nm laser might result from the fact that the 523 nm laser was generated from the so-called thermal coupled 3P­J and 1I6 multiplets and therefore an enlarged effective emission lifetime arising from the thermalization of the 3P1, 1I6 and 3P0 levels was in general attained. The large effective emission lifetime could lead to increased effective emission cross section for the green laser transition [17].

However, it

provides the highest Stokes efficiency and thus the lowest heat generation. Furthermore, it emits from a higher

level within the thermally coupled 3PJ and 1I6 multiplets

(see Fig. 1), which could lead to increased effective cross

sections at higher temperatures, making the laser even

benefit from moderate heating.

Less subject to thermal effect to some extent,

Can tolerance smaller intro-cavity mode size

a cavity length of ~0.3 m was also used to generate 522 nm self-mode-locked laser with a repetition rate of approximately 500MHz.

As a result, the total cavity length was ~1.8 m, which coincides well with the repetition rate measured to be ~85 MHz shown in Figure 5(c).

corresponding power

spectrum is shown in Fig. 3(c).

The laser was cw mode locked at ~85 MHz with only weak noise, and the difference between the peak of mode-locked frequency and that of relaxation oscillation frequency was experimentally found to be larger than 42 dBm.

might be restricted to great extent by

**4. Conclusion**

**Acknowledgments**

The authors wish to acknowledge the financial support from the National Natural Science Foundation of China (61275050, 61605069), National key Research and Development Program of China (2016YFB0701002).

**References**

[Novel self-mode-locking mechanism in narrow-band lasers]. Y. Bai, S. Chen, Z. Wang, and G. Zhang, Appl. Phys. Lett. 63, 2597 (1993).

[SESAM mode locked red praseodymium laser]. M. Gaponenko, P. W. Metz, A. Härkönen, A. Heuer, T. Leinonen, M. Guina, T. Südmeyer, G. Huber, and C. Kränkel, “SESAM modelocked red praseodymium laser,” Opt. Lett. 39, 6939–6941 (2014).